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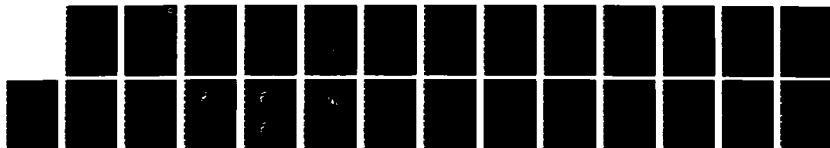
PROTON BEAM DEPOSITION HEATING AND RADIATION FROM AN
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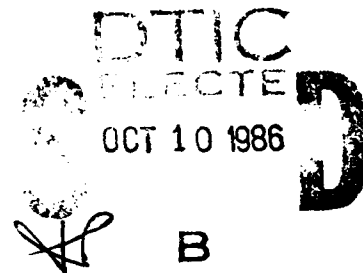


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Proton Beam Deposition, Heating, and Radiation from an Iron Plasma

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*Plasma Radiation Branch
Plasma Physics Division*



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PROTON BEAM DEPOSITION, HEATING, AND RADIATION FROM AN IRON PLASMA

I. Introduction

With current charged particle beam technologies, it is possible to bombard targets with intense proton beams. The interest in charged particle beam-target interactions is enhanced by the expectation that the beam-target energy coupling can be understood and modelled, to first approximation, in terms of well understood classical, linear processes, as opposed to the non-linear effects involved in laser-target coupling.¹ Thus, much effort is directed toward investigating these phenomena. Significant applications of this research include the use of ion beams for x-ray generation, inertial confinement fusion, and weapons lethality and vulnerability studies.

II. Theoretical Model

The interaction of a monoenergetic one MeV proton beam with a planar aluminum slab of 15 μm thickness has been previously investigated.² This beam was assumed to consist of a square-shaped pulse of 10 nanosecond duration with a flux of 10^{26} protons $\text{cm}^{-2} \text{sec}^{-1}$. The beam intensity on target was $1.6 \times 10^{13} \text{ W/cm}^2$, which is comparable to intensities available with current pulsed-power generators.

In the present study, the interaction of this same proton beam with a planar 8.0- μm -thick iron slab is treated. This distance approximately corresponds to the range of a one MeV proton in cold solid density Fe^3 .

The scheme for treating the proton beam-iron slab interaction is fundamentally the same as that described in Ref. 2 for treating the proton beam-Al slab interaction. The hydrodynamic response is treated via a one-dimensional model with a sliding-zone version of flux corrected transport. A special gridding algorithm moves zones in a Lagrangian fashion and adjusts the mesh in order to resolve steep gradients in the flow. A single temperature model is employed, since equilibration times are generally short compared with hydrodynamic timescales in this problem. A collisional radiative equilibrium treatment is used to obtain the atomic

level populations. Radiation transport is treated with a hybrid scheme which uses multifrequency ray tracing for the continua, and a probability-of-escape method for lines. The radiation transport is coupled self-consistently into the ionization dynamics. The proton stopping power includes contributions from bound atomic electrons, free plasma electrons, and plasma ions. These features are discussed in detail in Ref. 2.

The Fe atomic model used in this calculation did not have excited level structure in ionization stages below Fe XVI. Only ground states were carried for stages I through XV. While this lack of level structure and the resulting omission of lines from the radiation transport scheme can modify the energy balance in regimes where these lower stages are highly populated, it is estimated that the effect is less than a factor of two on the energy lost due to radiation.

A major difference between this work with Fe as the target material and the previous work with Al as the target material is in the treatment of the bound electron stopping power. For the proton-Al^{+q} interaction, where q is the ionic charge, detailed stopping power cross sections are available;⁴ for the proton-Fe^{+q} interaction, no such data are currently available. In the present treatment, stopping power due to bound electrons is calculated from a method proposed by Mehlhorn.⁵ In this approach, stopping power is calculated from the Bethe equation (which is valid at high projectile energies) and from the LSS model (which is valid in the low projectile energy regime); then the stopping cross section is taken to be the lesser of these two results; i.e.,

$$S_b(E) = \min [S_{\text{Bethe}}(E), S_{\text{LSS}}(E)] .$$

Mehlhorn⁵ has found that, by including shell corrections in the Bethe equation, he can adequately represent cold target stopping power results. By applying this model to cold solid density Fe, it has been verified that the Andersen-Ziegler³ results for proton stopping power can be reproduced very well.

The Bethe model for the stopping cross section is given by⁵

$$S_{\text{Bethe}}(E) = \frac{4\pi Z_B^2 (Z-\bar{Z}) e^4}{mV^2} \left\{ \ln \Lambda - \beta^2 - \sum_{i=1}^4 b_i [\log(E)]^i \right\} ,$$

$$\text{where } \Lambda = \frac{2\pi m V^2}{\langle I \rangle} \gamma^2 .$$

Here Z_B is the effective charge of the projectile ion, Z is the atomic number of the target, \bar{Z} is the average charge of the target ions, e and m are the electron charge and rest mass, E and V are the projectile energy and velocity, $\beta = V/c$, where c is the vacuum velocity of light, $\langle I \rangle$ is the average ionization potential of the target ions, and the b_i are least-squares fit parameters derived by Andersen and Ziegler³ to calculate shell corrections to the Bethe stopping power.

$$\gamma = (1 - \beta^2)^{-1/2}.$$

The LSS model is described by the following equations:⁵

$$S_{LSS}(E) = \left\{ \frac{0.0793 Z_B^{2/3} (Z - \bar{Z})^{1/2} (1+A)^{3/2}}{(Z_B^{2/3} + (Z - \bar{Z})^{2/3})^{3/4} A_2^{1/2}} \right\} \frac{\sqrt{E_L}}{R_L} \sqrt{E}$$

$$\text{where } E_L = \frac{Z_B(Z - \bar{Z})(1+A)e^2}{0.4683 \times 10^{-8} A} \times [Z_B^{2/3} + (Z - \bar{Z})^{2/3}]^{1/2},$$

$$R_L = \frac{(1+A)^2}{4\pi A} \frac{[Z_B^{2/3} + (Z - \bar{Z})^{2/3}]}{0.219 \times 10^{-16}}.$$

$A = A_2/A_1$, where A_2 is the atomic weight of the target ion, and A_1 is the atomic weight of the projectile ion.

The effective charge of the projectile ion is assumed to be given by the Brown and Moak⁶ relation

$$Z_B/Z_1 = 1 - 1.034 \exp[-V/(Z_1^{0.69} \times V_b)]$$

where Z_1 is the atomic number of the projectile ion and V_b is the Bohr velocity 2.188×10^8 cm/sec.

To calculate an average ionization potential for a target ion of charge q , a scheme proposed by Mehlhorn⁵ was adopted;

$$\langle I_q \rangle = (Z/Z_n)^2 \langle I_n \rangle,$$

where Z_n is the atomic number and $\langle I_n \rangle$ is the average ionization potential of a neutral atom with the same electronic configuration as the target ion of charge q .

Once the stopping cross section is obtained, the stopping power due to bound atomic electrons is given by

$$\left[\frac{dE}{dx} \right]_b = \frac{\rho(x)}{m_I} S_b(E)$$

where $\rho(x)$ is the mass density of target ions in the plasma at position x , and m_I is the mass of the target ion.

Stopping power due to free plasma electrons and plasma ions is calculated from the same models as used in the previous treatment of an Al target.² Enhanced stopping power due to collective beam effects is included, although in the case treated here, this effect is small.

III. Results

Some results of the proton beam-Fe slab target interaction simulation will now be given. Figure 1 shows the density, temperature, and average charge variations at the front and rear surfaces of the plasma as the interaction evolves. At the front surface, temperature and \bar{Z} rise very quickly and attain approximate steady state values within the first few nanoseconds. The density drops rapidly at first, then drops more slowly at later times, but it is always decreasing as the plasma expands. There is some slight cooling and a noticeable drop in \bar{Z} after the beam is shut off at 10 nsec. The temperature and \bar{Z} rise at the rear surface is much more gradual and does not begin to level off until near the end of the beam pulse. The density decreases here also as the plasma expands. The peaks in the rear surface plots at early times are due to the arrival of a shock wave.

Density and temperature profiles at different times in the plasma evolution are shown in Fig. 2. In these plots, the beam is incident from the right. The expansion of the plasma with time is easily discerned. At early times, there are significant gradients near the rear surface. As time goes on, these variations disappear and the plasma becomes more

uniform, especially at about 13 nsec, three nsec after the beam has been turned off. After three nsec, the temperature throughout most of the plasma remains fairly steady. At 1.7 nsec, the beam is stopped in the plasma interior and does not reach the backside to heat it up. This is also true at later times; the temperature increase at the rear surface arises from heating due to radiation transport from the interior of the plasma.

Calculated front and rear spectra at these times are displayed in Fig. 3. The lack of level structure in the lower ionization stages is demonstrated by the paucity of lines in Fig. 3a and in the rear spectra in Figs. 3b and 3c. Lines appear in the spectra when the plasma gets hot enough so that ions with level structure have a significant abundance. The peaks in the rear spectra, except at the lowest photon energies, arise from radiation from the plasma interior "shining through" opacity "windows". The dominant contribution to the front surface spectra generally comes from the outermost zone, although "shine-through" contributions from the interior can produce some peaks in the curve. As can be noted from Figs. 1 and 2, the plasma is nearly uniform at 13.15 nsec; hence the front and rear spectra are very similar. The front spectra becomes more intense at the higher frequencies up to 10 nsec, but, as might be expected, the 13.15 nsec front spectra is less intense than at 10 nsec due to plasma cooling after beam shutoff.

Figure 4 gives the energy history and partitioning of the beam-target interaction; "dep" is the energy deposited by the beam, and "sum" is the total energy remaining in the plasma in various forms. After the beam shutoff at 10 nsec, the total, internal, and potential energies decrease as the plasma cools; only the kinetic energy continues to rise as the plasma expands. At early times, the energy lost by radiation is insignificant, but by 13 nsec, this represents about a third of the deposited energy. Thus, radiation represents a significant energy loss. If the Fe atomic model had included level structure and line radiation in the lower ionization stages, the radiation loss would be greater, particularly at early times. However, as stated earlier, this lack of line structure is presumed to have an effect of less than a factor of two on the radiation loss.

IV. Summary

The interaction of a one-MeV proton beam with an intensity of 16 terawatts/cm² with a planar 8.0- μ m-thick iron slab has been simulated using a fully self-consistent, one-dimensional hydrodynamic-ionization-radiation transport code. An extensive atomic-level-structure model of Fe was employed, and atomic populations as functions of temperature and density were calculated from a collisional-radiative-equilibrium assumption. Radiation transport was treated with a hybrid scheme that combines a multi-frequency formalism for the continuum with a probability-of-escape method for the lines.

Profiles of the plasma density and temperature at various times in the plasma evolution, and corresponding front and rear emission spectra at these times, have been presented. Front and rear surface density, temperature, and average ionic charge has been discussed, and the energy partitioning among the various degrees of freedom has been shown.

Due to limitations of computer storage and other factors, detailed atomic structure was only included for Fe XVI and higher stages. Since the maximum average charge obtained was about 18, radiation losses are underestimated throughout the simulation; however, it is estimated that the error is less than a factor of two.

At these beam intensity levels, no K-shell emission was predicted, in contrast to the earlier study of the interaction of the same proton beam with an aluminum target.² Many L-shell lines are present during the latter part of the simulation.

There has been much interest recently in lasing schemes involving $2p^5 3p-2p^5 3s$ transitions in neonlike systems. One such transition is the $2p^5 3p(\frac{1}{2}, \frac{1}{2})_0-2p^5 3s(\frac{1}{2}, \frac{1}{2})_1$ transition at 254.9Å in Fe XVII. One of the reasons for undertaking this study is to see if this transition can be excited by proton beams. The atomic model did not have detailed jj levels for Fe XVII; the 3p levels were combined into two groups according to the spin of the $2p^5(^2P)$ inner shell. Nevertheless, it should be possible to make some estimates of the feasibility of inducing this transition. After about 5.5 nsec, these 3p levels had fractional population densities of $\sim 5-6 \times 10^{-3}$ to $\sim 10^{-2}$. During this time, the ion density was of the order 10^{20} ,

and the electron density of the order 10^{21} . The temperatures in the plasma during these times averaged about 180 eV, with some hotter zones rising above 200 eV. Under these conditions, it should be possible to see some lasing in this line.

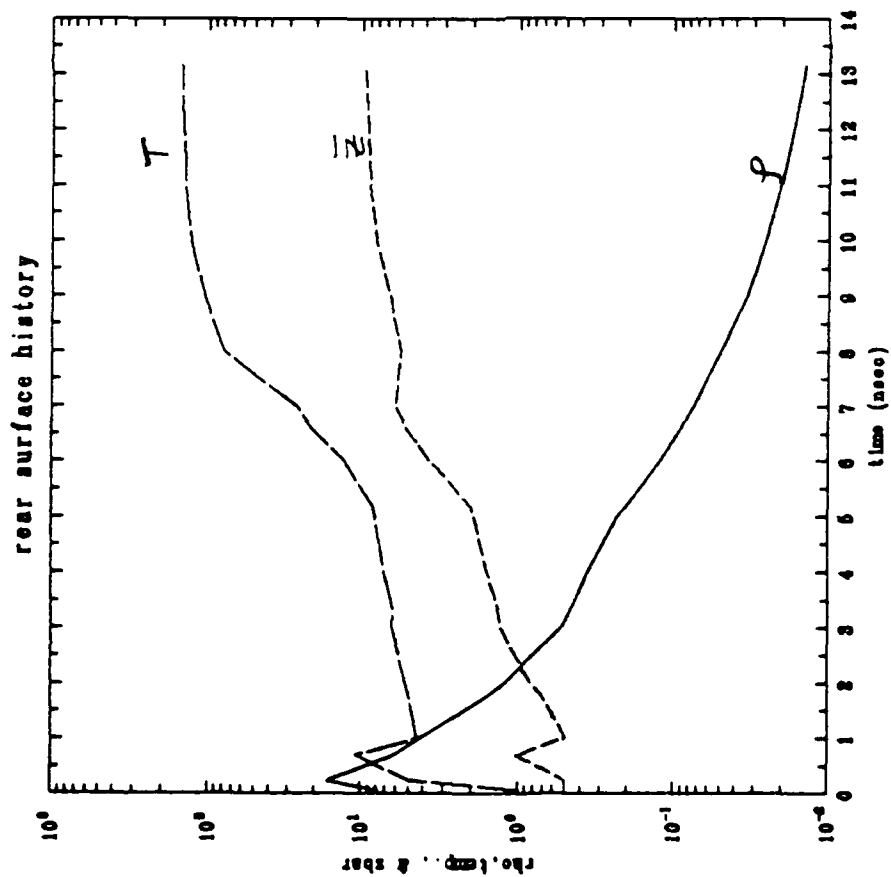
This treatment is one-dimensional; hence, field effects and lateral energy flow have been omitted. These factors could alter profile shapes and magnitudes, and stopping power in the plasma, and thus affect the results presented here.

V. ACKNOWLEDGMENTS

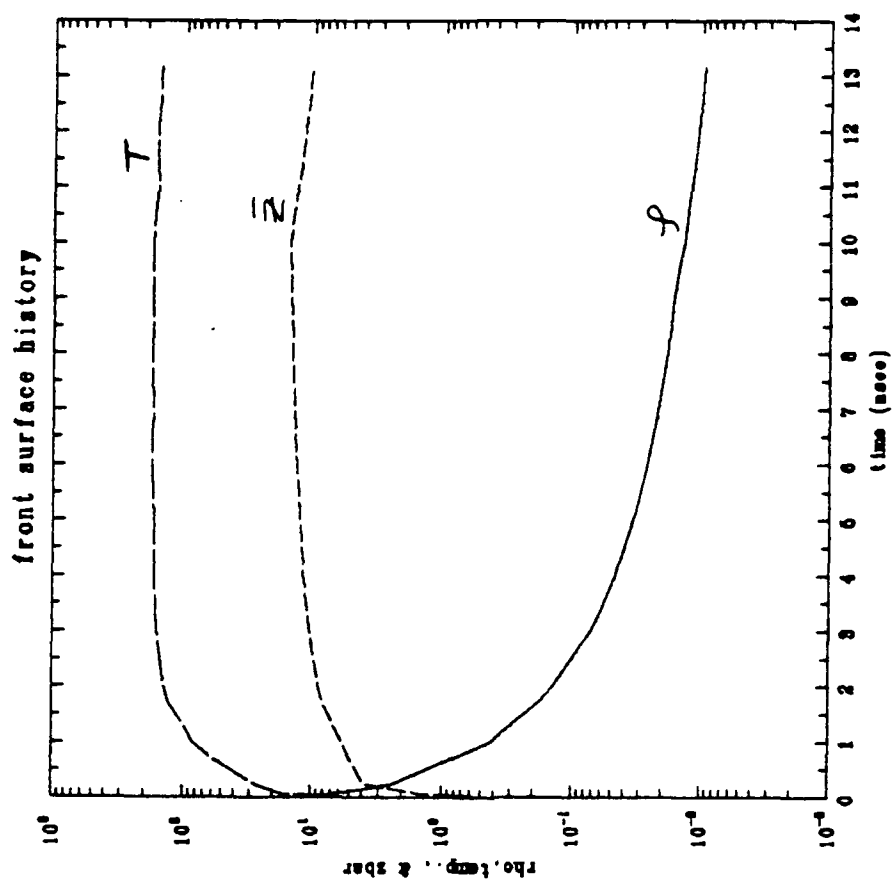
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(a)



(b)

Fig. 1 Temperature, density, and average charge variation at the front and rear surfaces with time.

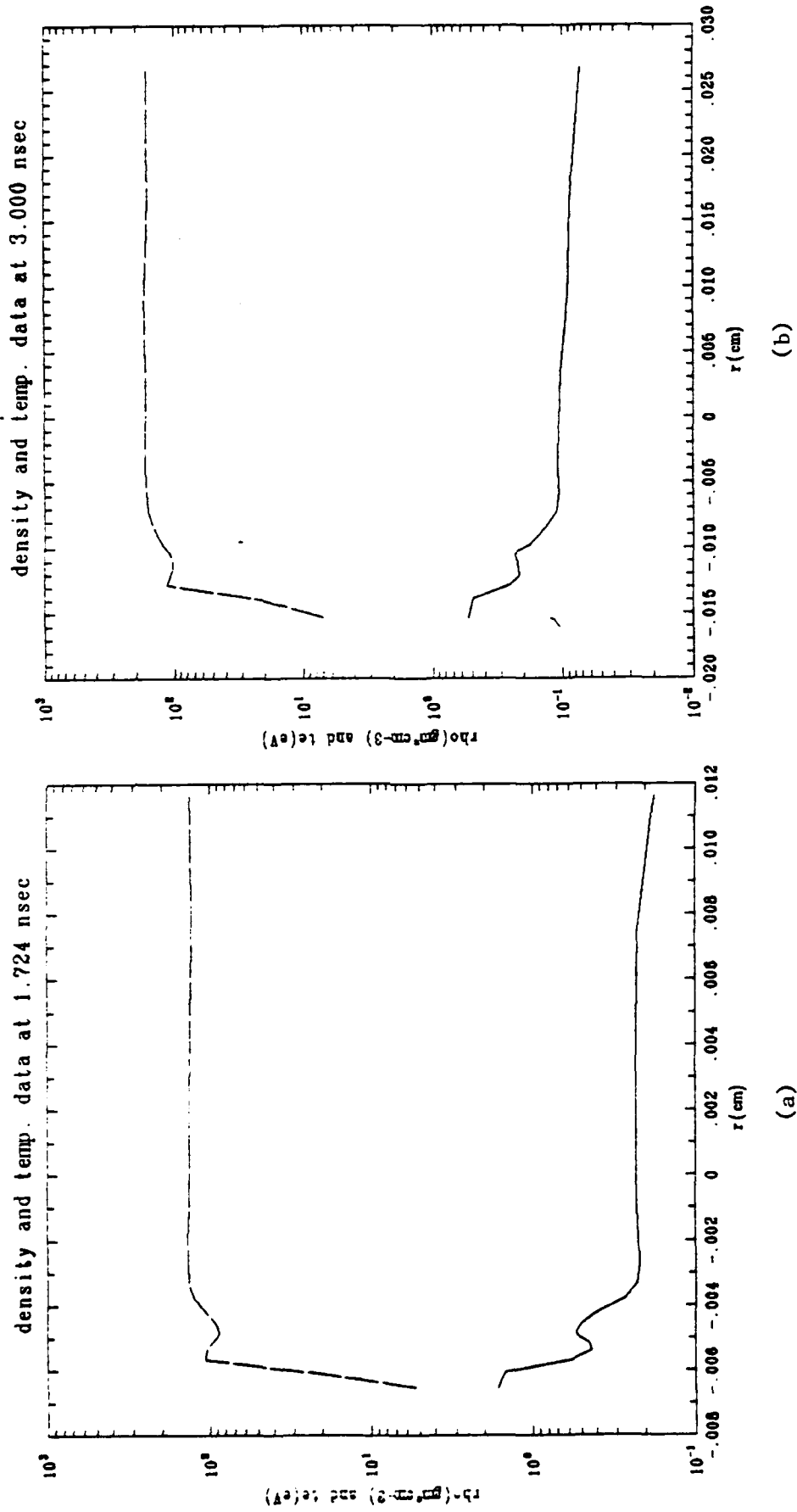
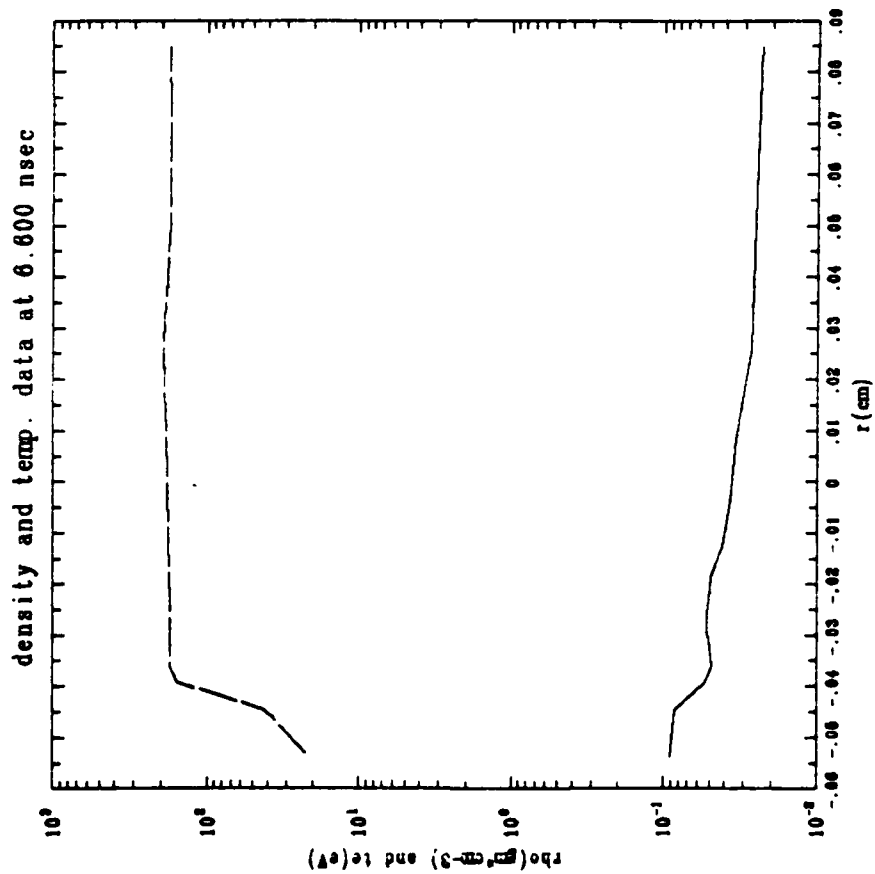
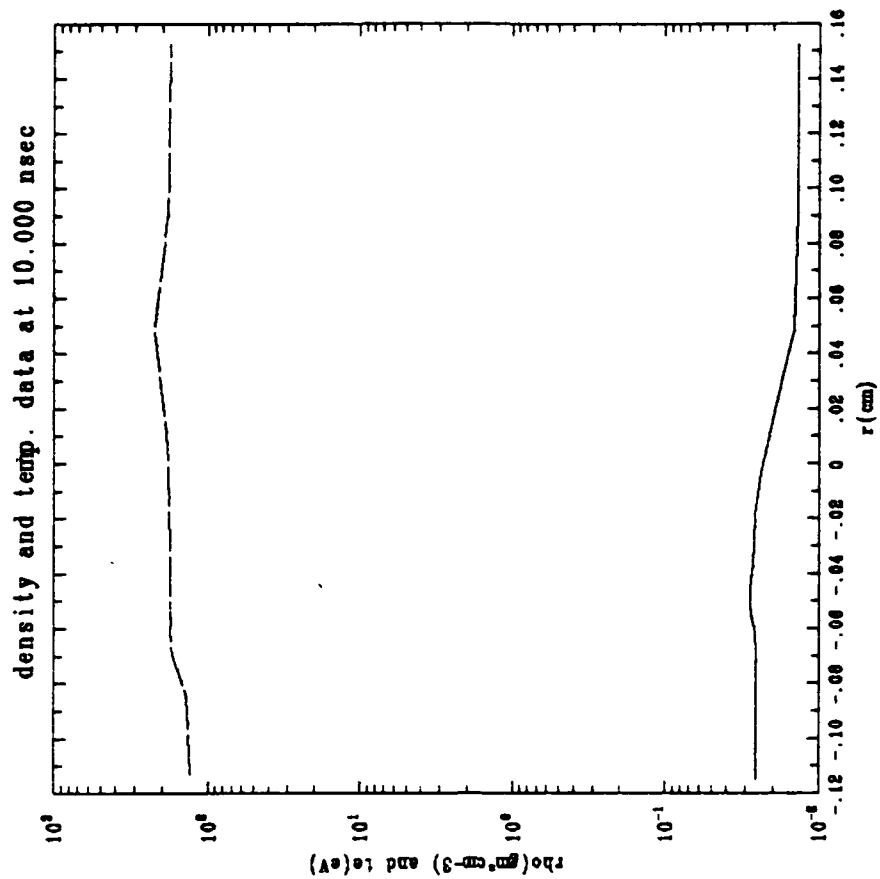


Fig. 2 Density and temperature profiles at different times in the plasma evolution.



(c)



(d)

Fig. 2 (Continued) Density and temperature profiles at different times in the plasma evolution.

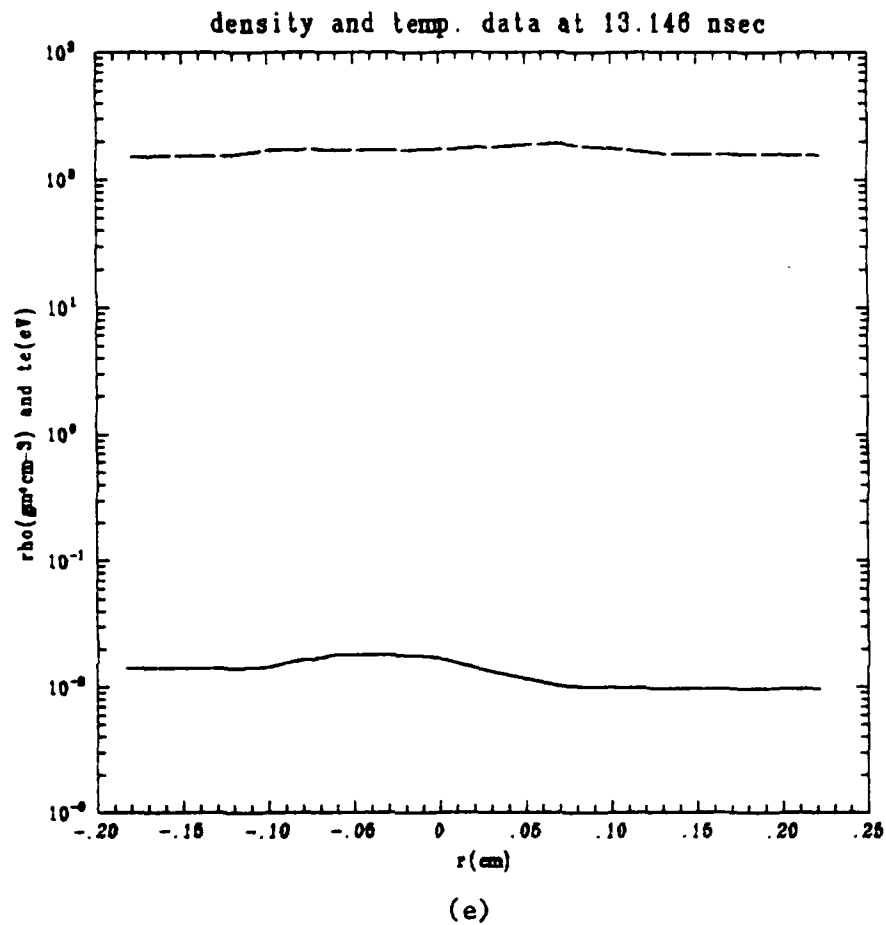


Fig. 2 (Continued) Density and temperature profiles at different times in the plasma evolution.

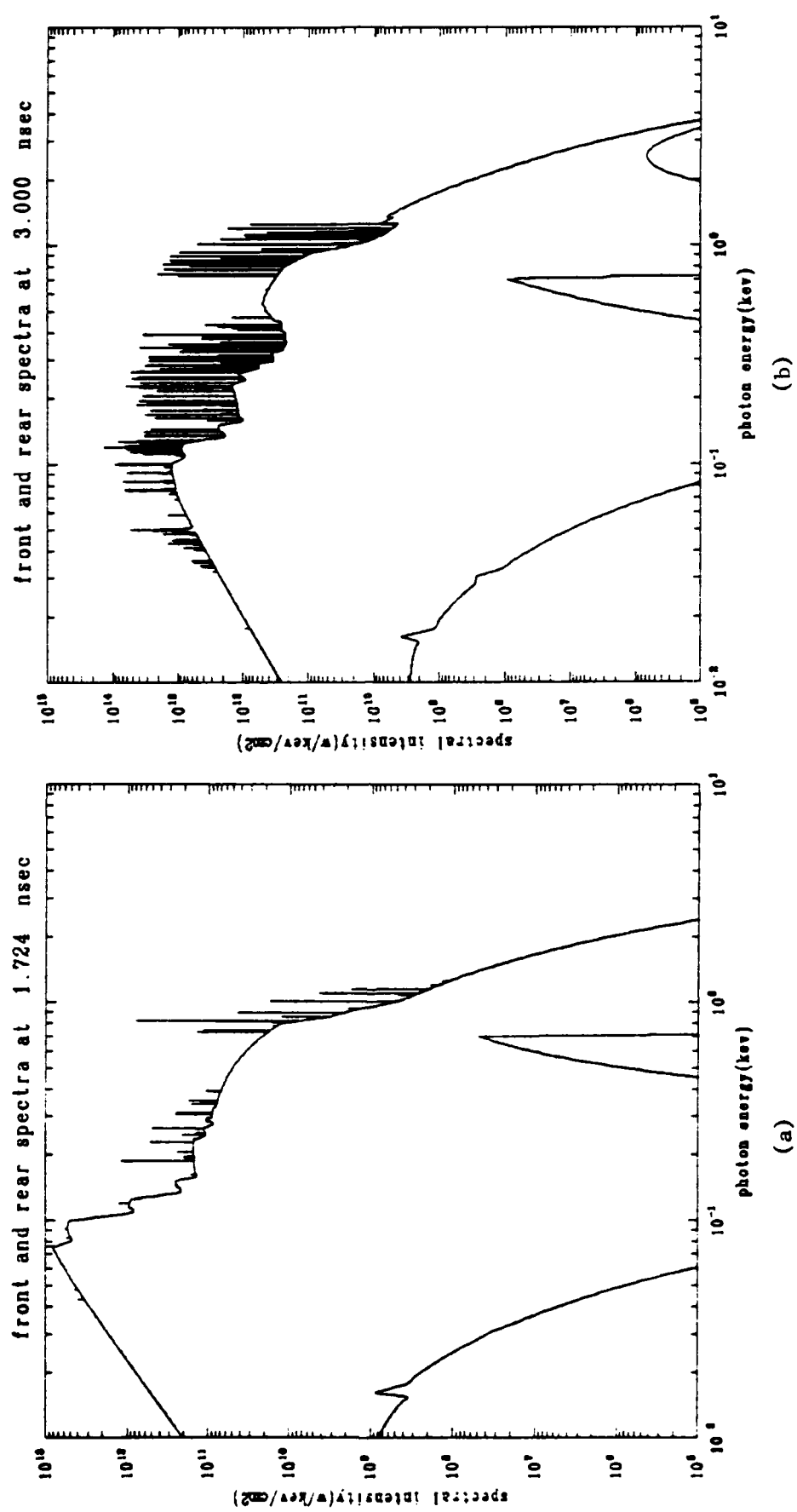
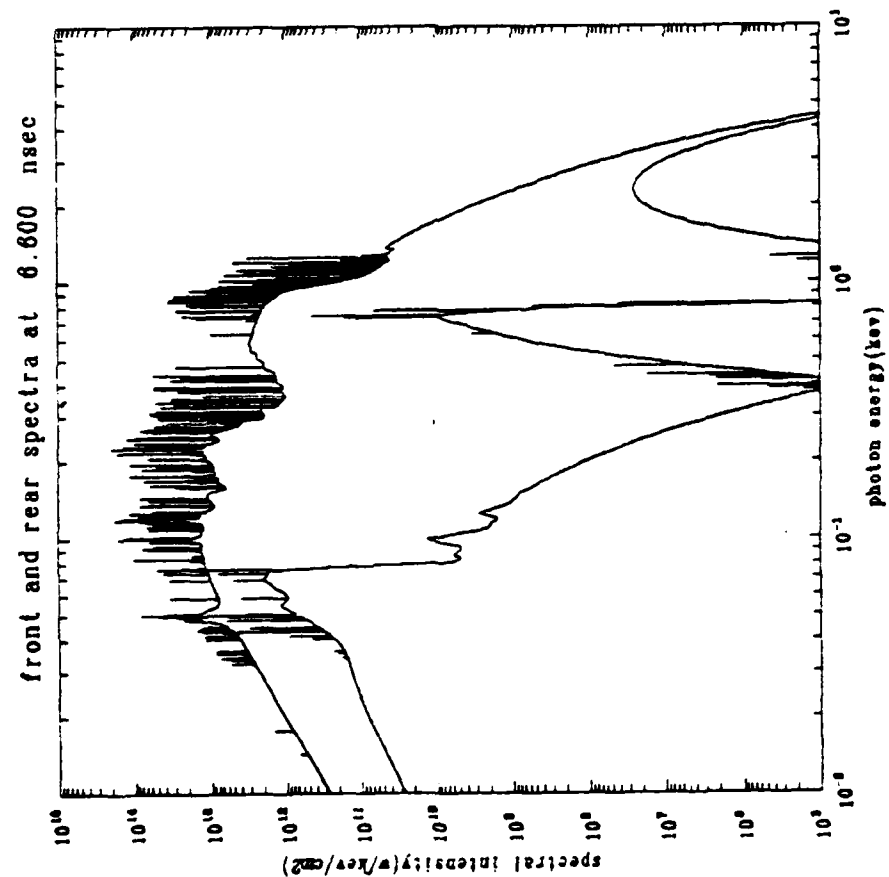
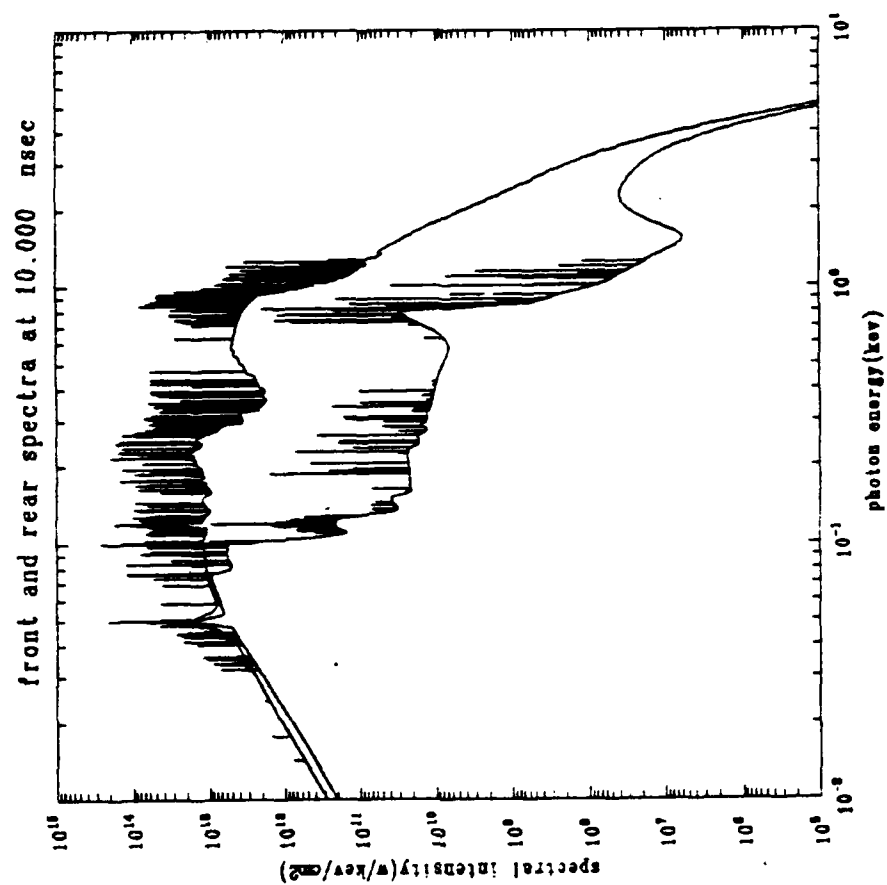


Fig. 3 Front and rear emission spectra at various times in the plasma evolution.

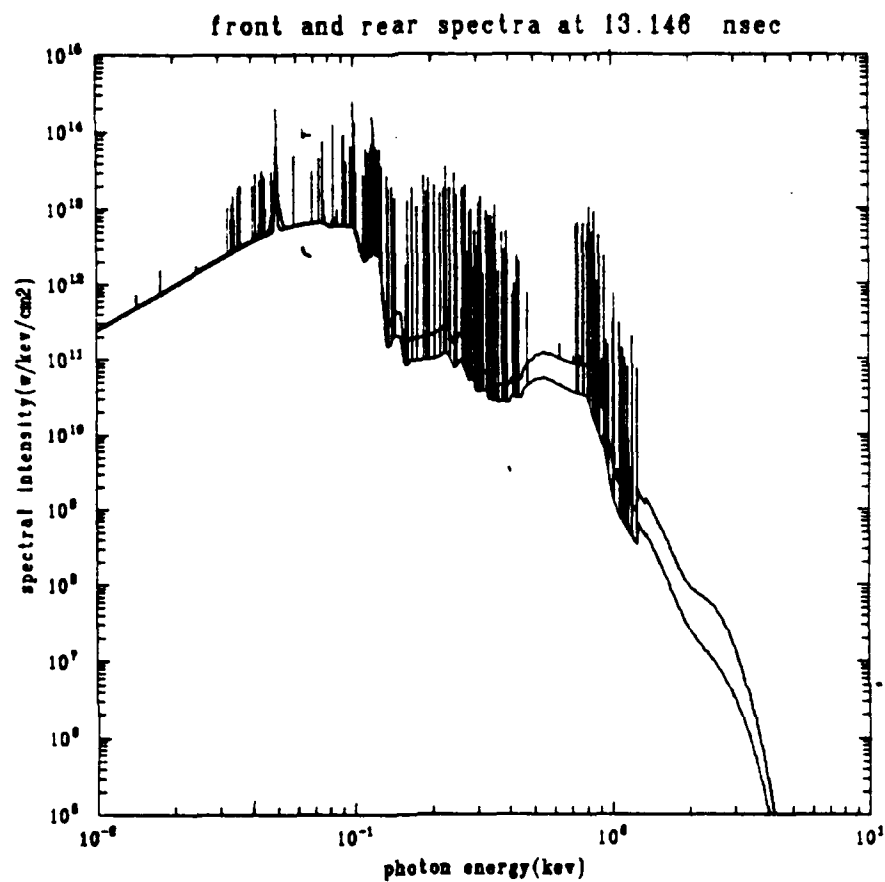


(c)



(d)

Fig. 3 (Continued) Front and rear emission spectra at various times in the plasma evolution.



(e)

Fig. 3 (Continued) Front and rear emission spectra at various times in the plasma evolution.

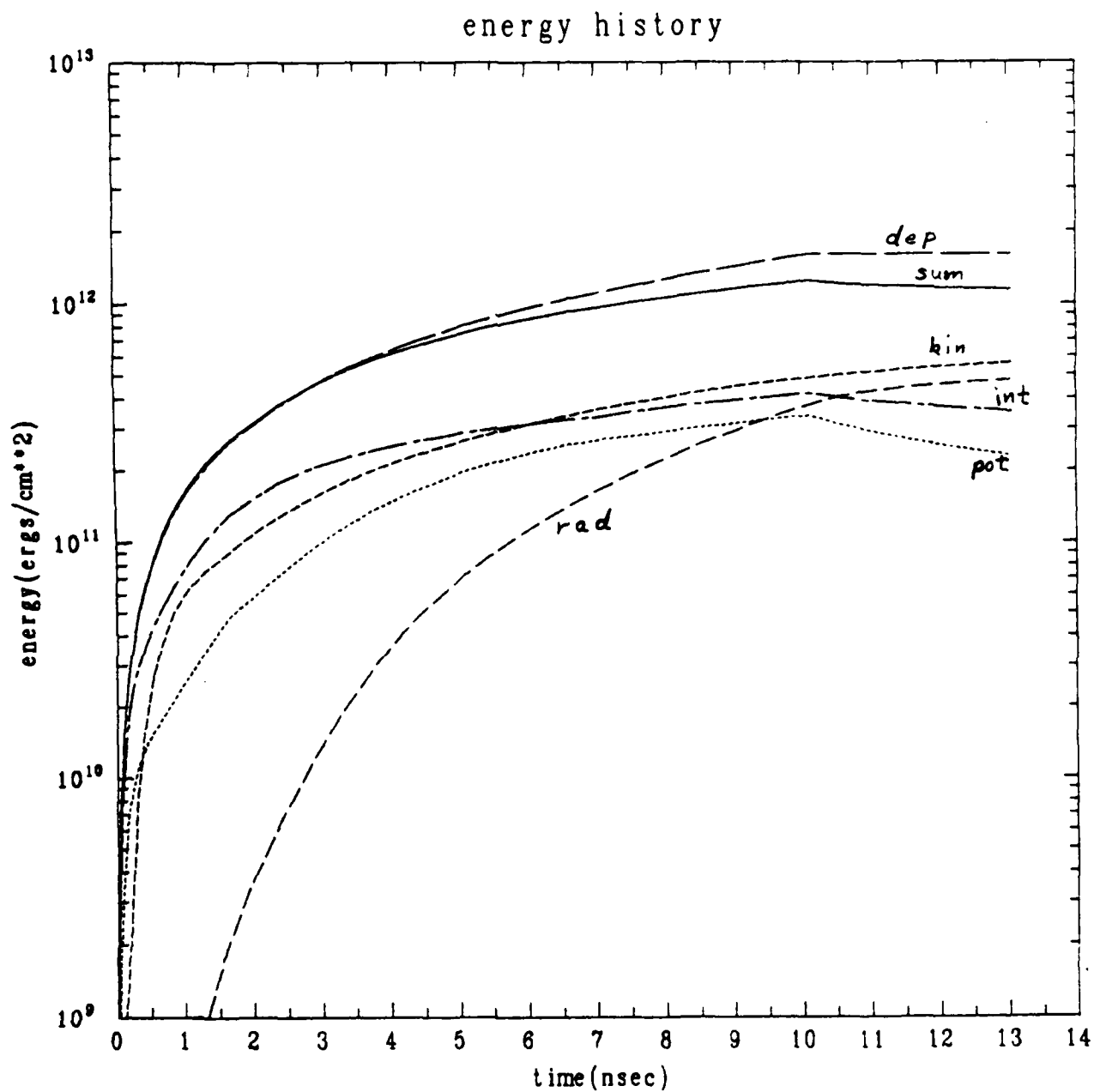


Fig. 4 Energy history of the beam-plasma evolution.

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